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<https://journals.aps.org/prb/abstract/10.1103/PhysRevB.107.115140>

To cite this version

Min Long, Zhen Zhan, Pierre A. Pantaleón, Jose Ángel Silva-Guillén, Francisco Guinea, and Shengjun Yuan. Electronic properties of twisted bilayer graphene suspended and encapsulated with hexagonal boron nitride (2023).

<https://repositorio.imdeananociencia.org/handle/20.500.12614/3284>

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Electronic properties of twisted bilayer graphene suspended and encapsulated with hexagonal boron nitride

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(Dated: January 25, 2024)

The recent observed anomalous Hall effect in magic angle twisted bilayer graphene (TBG) aligned to hexagonal boron nitride (hBN) and unconventional ferroelectricity in Bernal bilayer graphene sandwiched by hBN present a new platform to tune the correlated properties in graphene systems. In these graphene-based moiré superlattices, the aligned hBN substrate plays an important role. In this paper, we analyze the effects of hBN substrate on the band structure of the TBG. By means of an atomistic tight-binding model we calculate the electronic properties of TBG suspended and encapsulated with hBN. Interestingly, we found that the physical properties of TBG are extremely sensitive to the presence of hBN and they may be completely different if TBG is suspended or encapsulated. We quantify these differences by analysing their electronic properties and band topology. We found that the narrow bandwidth, band gap and local density of states are significantly modified by the aligned hBN substrates. Interestingly, these electronic properties can be used as a signature of the alignment in experiment. Moreover, the TBG/hBN superlattices in the presence or absence of the two-fold rotation symmetry response differently to the external electric field. For the TBG suspended in the hBN, application of an electric field results in the charge unevenly distributed between graphene layers, which can be used to tune the strength of the valley Hall effect or the anomalous Hall effect. Such rich topological phase diagram in these systems may be useful for experiments.

I. INTRODUCTION

Since the discovery of unconventional superconductivity and correlation effects in twisted bilayer graphene (TBG) near the magic angle by Cao *et al.* [1, 2], the so-called field of “twistronics” has become of great interest to the condensed matter community [3]. In this magic angle at approximately 1.1° [4], the system possesses flat bands near charge neutrality, which are responsible for most of these exotic behaviors. Interestingly, not only two graphene layers have been twisted, but also transition metal dichalcogenides [5, 6], monochalcogenides [7], hexagonal boron nitride (hBN) [8, 9] and black phosphorus [10], among others. Similar to the TBG, flat bands are observed in most of these twisted two-dimensional (2D) materials.

Experimentally, the devices are usually supported on a substrate which is typically a thin sample of hBN [11]. Although hBN has a large gap and has been thought to have a small impact on the electronic properties of the materials which are supporting, it is not the case. As we have shown recently [12] (see also [13–16]), hBN has an important effect on the electronic properties of TBG, when the sample is either supported or encapsulated. We have found that when the sample is placed on top of hBN

a gap opens due to the appearance of a mass term as a consequence of the breaking of \mathcal{C}_2 symmetry [12, 17]. Furthermore, a splitting in the bands appears due to layer degeneracy breaking. In fact, hBN affects the electronic properties of TBG even when the angle between TBG and hBN is far from alignment. When TBG is encapsulated between two layers of hBN the layer degeneracy can be recovered for certain angles, while the gap still appears. In a typical experimental setup, an electric field is usually applied to the samples in order to change their doping which could also modulate the electric structures of the samples [18, 19]. Therefore, the study of the electronic such as the local density of states (LDOS) of TBG in combination with a substrate and when an electric field is applied is very compelling.

In this work, we extend our previous study to investigate the aforementioned properties in both the case of TBG supported and encapsulated between hBN layers. In Sec. II we describe the atomic structure of our system and the methods that we employ to perform our calculations. Then, in Sec. III the layer degeneracy in electronic properties like LDOS and charge density distribution are studied. In Sec. IV we explore the response of these systems to an external field. In Sec. V we study the effect of encapsulating TBG on its band topology. Finally, we give a summary and discussion of our work.

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II. ATOMIC STRUCTURES AND NUMERICAL METHODS

A. The Atomic structures

In this work, we mainly focus on two structures. The first one is a trilayer system composed of TBG and a hBN layer lying on the bottom (TBG/hBN). The second one is a sandwich-like system where TBG is encapsulated by two hBN layers (hBN/TBG/hBN). TBG and hBN can be stacked in different ways to have commensurate structures [20]. In our case, we start by constructing an unrotated trilayer structure, where graphene and hBN bonds are parallel and the center of the supercell has an AAA stacking with a carbon site of graphene and a nitrogen (N) site of hBN sharing the same in-plane position, $(x, y) = (0, 0)$. Then we rotate the top layer graphene (G_{top}) and bottom layer hBN (hBN_{bot}) with angles θ_{tbg} and θ_{bot} with respect to the bottom layer graphene (G_{bot}), respectively. In our notation, positive angles correspond to counterclockwise rotations.

We define the lattice vectors of a hexagonal lattice as $a_1 = a(\sqrt{3}/2, 1/2)$ and $a_2 = a(\sqrt{3}/2, -1/2)$ being a the lattice constant. For graphene $a_g = 0.246$ nm, and for hBN $a_{hBN} = 0.2503$ nm. The lattice mismatch between graphene and hBN is $\delta \approx 1.8\%$. The twist angle of TBG can be solely determined by a coprime integer pair (m, n)

$$\theta_{tbg} = 2 \arcsin \frac{(m-n)}{2\sqrt{m^2 + mn + n^2}}, \quad (1)$$

with moiré length

$$L_{tbg} = \frac{a_g(n-m)}{2|\sin(\theta_{tbg}/2)|} = a_g \sqrt{m^2 + mn + n^2}. \quad (2)$$

The moiré length of the G_{bot}/hBN_{bot} superlattice is [11]

$$L_{hBN} = \frac{(1+\delta)a_g}{\sqrt{\delta^2 + 2(1+\delta)(1-\cos\theta_{bot})}}. \quad (3)$$

A commensurate structure of TBG/hBN is obtained when

$$L_{tbg/hBN} = L_{tbg} = qL_{hBN} \quad (4)$$

where q is an integer. We focus on a system where $\theta_{tbg} = 1.05^\circ$ and $\theta_{bot} = 0.53^\circ$. For this angle combination, the periodicity of the moiré pattern of TBG is identical to that constructed by hBN and G_{bot} , and therefore a single moiré unit cell can be defined for the combined system [17, 21]. That is, the three moiré lengths have the same value $L_{tbg/hBN} = L_{tbg} = L_{hBN} = 13.4$ nm, and $q = 1$. To construct the hBN/TBG/hBN structure, we just add a second hBN layer (hBN_{top}) on the top of the trilayer structure, the twist angle between hBN_{top} and G_{bot} is also $\theta_{top} = 0.53^\circ$. It is important to note that, in order to keep the periodicity of these structures, the lattice constant of hBN is slightly modified.

In our case, this implies a strain of about 0.14%. Such a small value of strain will not affect the structural or electronic properties of hBN, which, therefore, will not change the properties of the TBG/hBN systems that we study in this work [12]. The schematics of TBG/hBN and hBN/TBG/hBN are shown in Fig. 1(a) and Fig. 2(a), respectively.

B. The tight-binding model

We adopt a combination of semi-classical molecular dynamics and a tight-binding (TB) model to investigate the electronic properties of TBG/hBN and hBN/TBG/hBN structures. After constructing a commensurate supercell, we use semi-classical molecular dynamics, which is implemented in LAMMPS [22], to fully (both in-plane and out-of-plane) relax the graphene layers in the two systems. For intralayer interaction between graphene, we use the reactive empirical bond order potential (REBO) [23]. For interlayer interaction between graphene, we use the registry-dependent Kolmogorov-Crespi (RDKC) potential developed for graphite [24]. We use the same RDKC potential for the C-B and C-N interlayer interaction, but with different strengths. The interaction strength of C-B and C-N are 60% and 200% with respect to the original C-C interaction, respectively [25]. The hBN layers are fixed in a flat configuration to mimic a bulk or a few layers of substrate. We assume that the relaxed structures keep the periodicity of the rigid cases.

The full TB Hamiltonian for graphene and hBN heterostructure can be written as [12]

$$\hat{H} = - \sum_{i,j} t(\mathbf{R}_i - \mathbf{R}_j) |\mathbf{R}_i\rangle \langle \mathbf{R}_j| + \sum_i \epsilon(\mathbf{R}_i) |\mathbf{R}_i\rangle \langle \mathbf{R}_i| + \sum_i V_D(\mathbf{R}_i) |\mathbf{R}_i\rangle \langle \mathbf{R}_i|, \quad (5)$$

where \mathbf{R}_i and $|\mathbf{R}_i\rangle$ represent the atom position and the atomic state at site i , respectively, $t(\mathbf{R}_i - \mathbf{R}_j)$ is the transfer integral between the atomic states at sites i and j , $\epsilon(\mathbf{R}_i)$ encodes the carbon, boron and nitrogen onsite energies and $V_D(\mathbf{R}_i)$ is the deformation potential resulting from the structural relaxation. For the onsite energy of boron, nitrogen and carbon atoms, we assume [26]:

$$\epsilon_B = 3.34 \text{ eV}, \quad \epsilon_N = -1.40 \text{ eV}, \quad \epsilon_C = 0 \text{ eV} \quad (6)$$

The lattice deformation leads to the emergence of periodic scalar and gauge potentials [27–33]. All these effects can be accurately considered in our TB model. To incorporate the relaxation effect into Hamiltonian in Eq. (5), we introduce the deformation potential term as [25]:

$$V_D(\mathbf{R}_i) = g_1 \frac{S(\mathbf{R}_i) - S_0}{S_0}, \quad (7)$$

where the screened deformation potential $g_1 = 4$ eV [34], $S(\mathbf{R}_i)$ is the effective area of site i that is modulated

by local deformations, and $S_0 = \sqrt{3}a/4$ is the effective area in equilibrium. For the transfer integral, we simply adopt the common Slater-Koster-type function for any combination of atomic species [26]:

$$-t(\mathbf{R}) = V_{pp\pi} \left[1 - \left(\frac{\mathbf{R} \cdot \mathbf{e}_z}{R} \right)^2 \right] + V_{pp\sigma} \left(\frac{\mathbf{R} \cdot \mathbf{e}_z}{R} \right)^2, \quad (8)$$

where

$$V_{pp\pi} = V_{pp\pi}^0 e^{-\frac{R-a_0}{r_0}}, \quad (9)$$

$$V_{pp\sigma} = V_{pp\sigma}^0 e^{-\frac{R-d_0}{r_0}}. \quad (10)$$

In the above equation, \mathbf{e}_z is the unit vector perpendicular to the graphene plane, $R = |\mathbf{R}|$, $a_0 = a_g/\sqrt{3} \approx 0.142$ nm is the C-C distance, $d_0 = 0.335$ nm is the interlayer distance, $V_{pp\pi}^0$ and $V_{pp\sigma}^0$ are the intralayer and interlayer transfer integrals between nearest neighbor atoms, respectively. We take $V_{pp\pi}^0 \approx -2.7$ eV, $V_{pp\sigma}^0 \approx 0.48$ eV. The parameter r_0 is the decay length of the transfer integral and is chosen as $0.184a_g$ so that the next nearest intralayer coupling becomes $0.1V_{pp\pi}^0$. For atoms whose distance is more than 0.6 nm, we set $t(\mathbf{R}) = 0$ since for larger distances the value of hopping energy is small enough to be safely neglected.

C. The electronic properties

Once the TB Hamiltonian is constructed, we can calculate the electronic properties of the TBG/hBN superlattices. Since the TBG/hBN and hBN/TBG/hBN structures contain tens of thousands of atoms, we use a tight-binding propagation method (TBPM) to obtain the density of states (DOS). The TBPM is based on the numerical solution of the time-dependent Schrödinger equation and requires no diagonalization processes, which is implemented in our home-made TBPLaS simulator [35–37]. In TBPM, a random initial state $|\phi_0\rangle$ is used with $\langle\phi_0|\phi_0\rangle = 1$. The density of states is calculated as a Fourier transform of the time-dependent correlation function

$$D(E) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{iE\tau} \langle\phi_0|e^{-iH\tau/\hbar}|\phi_0\rangle d\tau \quad (11)$$

In the TBPM, the convergence can be guaranteed by averaging over different initial states $|\phi_0\rangle$. For large enough systems, the results are converged with only one random initial state. The LDOS is obtained via the recursion method in real space based on the Lanczos algorithm [38]. The eigenstates and eigenvalues are obtained by direct diagonalization of the Hamiltonian in Eq. (5).

III. TBG HETEROSTRUCTURES

A. TBG supported on hBN

In free standing TBG, the layer degree of freedom is disentangled from spin and valley, forming eight-fold degeneracy in the low energy flat bands. As shown in Fig. 5, the conduction and valence flat bands are connected by Dirac points at the K and K' points of the moiré Brillouin zones (mBZ). These Dirac points are protected by $\mathcal{C}_2\mathcal{T}$ symmetry, where \mathcal{C}_2 is the two-fold rotation operator and \mathcal{T} is the time-reversal operator. If $\mathcal{C}_2\mathcal{T}$ is broken, the Dirac points will become gapped. As we discussed in our previous work [12], in a TBG/hBN superlattice, the hBN substrate introduces two contributions to the TBG system: the first one is the energy difference between nitrogen and boron atoms. This gives rise to different adhesion energies in the TBG/hBN system and breaks the \mathcal{C}_2 inversion symmetry, which is the main source of this mass gap. The second contribution is the lattice relaxation effects that give rise to a deformation potential and a pseudo-magnetic field [39]. Such relaxation effect ensures the persistence of a gap opening in the Dirac point for angles between TBG and hBN far from alignment [40].

The substrate effects can be clearly observed in the band structure shown in Fig. 1(b). Narrow bands are separated by a gap of around 30 meV due to the breaking of \mathcal{C}_2 symmetry [40–51]. This first gap (energy difference between the flat bands at K) is reduced when the angle θ_{bot} increased [12]. Moreover, the presence of a substrate acting on a single graphene layer breaks the mirror symmetry between layers and their Dirac cones are shifted in energy. This layer degeneracy breaking is responsible for the observed splitting between narrow bands. This splitting is more obvious in the remote bands located at around ± 0.1 eV. As we will discuss later, a perpendicular electric field will further increase this splitting. The narrow bands show a significant electron-hole asymmetry due to the strong superlattice potential induced by the hBN. Compared to TBG (see Fig. 5) the flat bands become dispersive and the peaks in the DOS of the TBG at charge neutrality are smoothed giving rise to an insulating structure in Fig. 1(b). In graphene/hBN superlattices, secondary Dirac cones appear at higher energies, which can be attributed to the moiré potential [25]. In the TBG/hBN structure that we study, the periodicity of the moiré pattern is 14 nm. This entails the secondary Dirac cones to be located at around 0.14 eV (outlined by a dashed rectangle in Fig. 1(b)). The states of the induced Dirac cones are mainly localized at the ABC or ACB stacking regions (the results not shown here), which is similar to the results of graphene/hBN superlattices [25].

The hBN substrate also significantly modifies the local properties of the TBG/hBN superlattice. In free standing TBG (see Fig. 5), the flat band states are mainly localized in both sublattices A and B around the AA

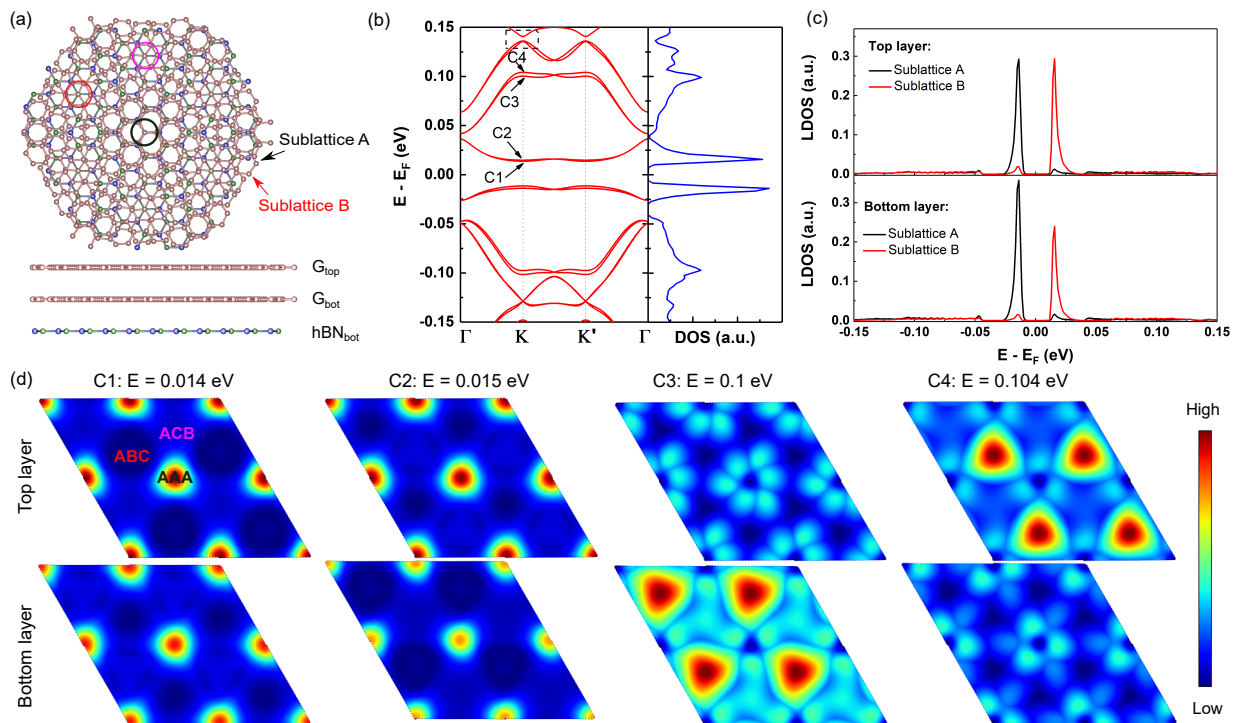


FIG. 1. (a) Top (upper panel) and side (lower panel) views of the trilayer structure. The center of the top view has an AAA high-symmetry stacking, where the carbon atom of both graphene layers and the nitrogen atom of hBN share the same in-plane position. Sublattices A and B of graphene are specified. The high-symmetry stackings of AAA, ABC and ACB are outlined with black, red and purple circles, respectively. (b) Band structure and density of states of states of the trilayer structure. The black arrows indicate various significant optical transitions. (c) Local density of states of sublattices A and B in the AAA stacking region. The LDOS of atoms from top (upper panel) and bottom (lower panel) graphene layers are plotted separately. (d) Eigenstates $|\psi|^2$ in real space. The eigenstates are calculated by diagonalizing the Hamiltonian in Eq. (5). The corresponding energies $C1$, $C2$ and $C3$ of these eigenstates are illustrated in (b).

stacking regions. If we look at the calculated LDOS of sublattices A and B at the AAA stacking region, as shown in Fig. 1(c), the peak in the sublattice B is lower than that of the sublattice A in the bottom layer. This is due to the role that hBN is playing as a substrate. The hBN substrate breaks the sublattice symmetry in the G_{bot} , making the LDOS peaks in the sublattice B lower than the sublattice A. On the contrary, the hBN substrate has negligible influence on the top layer. The states from G_{top} and G_{bot} contribute unequally to the conduction band and valence band, which is the natural result of the breaking of the layer degeneracy. In some works, the hBN substrate effect on the TBG is introduced via an effective periodic potential acting on the nearest graphene layer [26, 27]. Our results indicate again that this approach is also correct. We also plot the real space wave function of the narrow bands. Different from the TBG case shown in Fig. 5, where states in the narrow bands are localized around the AA centers. In the presence of hBN we observe some states localized in the ABC or ACB stacking regions. That is, a small part of states from the C1 band are localized in the ACB region (see labels in Fig. 1(b)), whereas some states from the C2 band are in the ABC region. Such difference is a consequence of the

substrate potential that redistribute the charges within the moiré unit cell. The states from a higher energy of C3 are mainly localized in the ACB region of the bottom layer.

B. TBG encapsulated in hBN

We now consider a tetralayer structure where, as shown in Fig. 2(a), TBG is encapsulated by two hBN layers. As described in our previous work [12] there are several possible stacking configurations for the tetralayer structure (see also [20]). In this work we choose the twist angle and stacking configuration such that the mirror symmetry between layers is recovered, this is given by the condition $|\theta_{bot}| = |\theta_{top}|$. Therefore, we could tune the layer degeneracy by adding or removing one of the hBN layers (as we will describe in the following section, an electric field can also be used to tune the layer degeneracy). It is important to mention that the breaking of the layer degeneracy is because the Dirac cones in each graphene layer are being affected by their nearest hBN layer. A direct consequence of recovering the layer degeneracy is the disappearance of the splitting in the flat bands, as

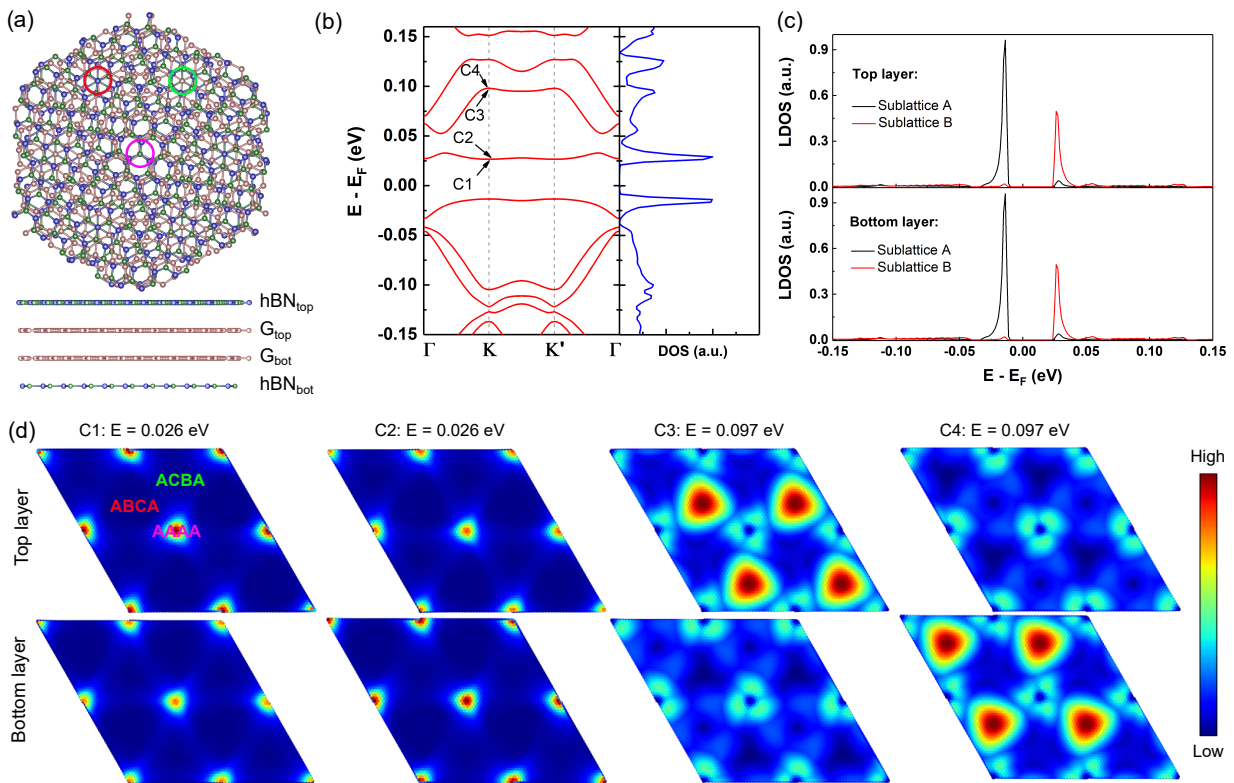


FIG. 2. (a) Top (upper panel) and side (lower panel) views of the tetralayer structure. The center of the top view is the AAAA high-symmetry stacking, which has the carbon atoms of graphene and nitrogen atoms of hBN share the same in-plane position. The high-symmetry stackings of AAAA, ABCA and ACBA are outlined by purple, red and green circles, respectively. (b) Band structure and density of states of tetralayer structure. The black arrows indicate various significant optical transitions. (c) Local density of states of sites in the AAAA stacking region. (d) Eigenstates $|\psi|^2$ in real space. The corresponding energies, C1, C2 and C3 of these eigenstates are illustrated in (b).

we can see from the comparison between the band structure in Fig. 1(b) with that in Fig. 2(b). The splitting does not only disappear in the narrow bands near the Fermi level, but also in the high energy bands, indicating an intrinsic change in the electronic structure of the encapsulated structure. Compared with TBG, both middle bands become narrower. The gap between the lower narrow bands has a value of around 50 meV. Moreover, as we can see from the LDOS calculation, Fig. 2(c), the states from G_{top} and G_{bot} contribute equally to both narrow bands. Another noteworthy result is that the states of both bands have different contributions from different sublattices. That is, both the top and bottom layers have a different sublattice charge distribution, with the LDOS peaks of the sublattice B lower than the sublattice A. Regarding the charge density maps in Fig. 2(d) and compared to the TBG/hBN superlattice, the states of the narrow bands become more localized in the AAAA regions, forming a triangular shape with smaller area. Moreover, the states of the two upper narrow bands have a quite similar localization in real space. For some other tetralayer structures, for example, structures without mirror symmetry as those discussed in our previous work [12], the splitting of the flat bands is pre-

served. Since the sublattice asymmetry is induced from the hBN substrate, in the tetralayer structures without mirror symmetry, we still could expect a different charge distribution of sublattices A and B in both top and bottom graphene layers. We have found that the hBN substrate enhances the localization of the flat band states in the AAA regions, an additional hBN substrate will increase such enhancement.

IV. THE EFFECT OF A PERPENDICULAR ELECTRIC FIELD

It has been shown that a perpendicular electric field in graphene/hBN superlattices allows the tunability of their physical properties [52], such as the case of Bernal bilayer graphene on hBN (BG/hBN) [53], ABC stacked graphene on hBN (TG/hBN) [54] or twisted double bilayer graphene (TDBG) [55], among others. Interestingly, the bandwidth of the narrow bands in these three systems is reduced by increasing the electric field [52, 55, 56]. However, in some other graphene systems, for instance, the twisted trilayer graphene with mirror symmetry, the bandwidth is increased by an external electric

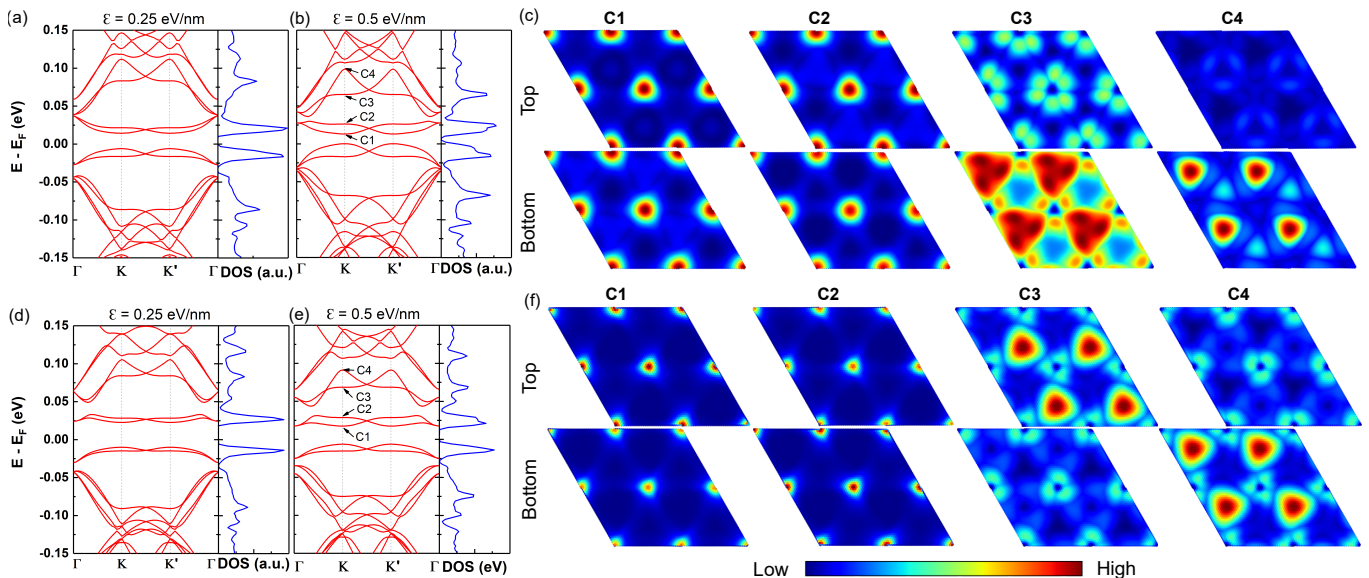


FIG. 3. (a)(b),(d)(e) Band structures and density of states for trilayer structure and tetralayer structure with different electric field, respectively. (c) and (f) Eigenstates, $|\psi|^2$, calculated from diagonalization of the Hamiltonian for the trilayer and tetralayer structure, respectively. The corresponding energies of these eigenstates are illustrated in (b) and (e). In the density plots, red and dark blue indicate the maximum and minimum values of the charge density, respectively. We note that the direction of the electric field with positive value \mathcal{E} has the same direction as the unit vector \mathbf{e}_z .

field [57, 58]. In TBG a perpendicular electric field shifts the degeneracy of the Dirac cones [59–61] and does not open a gap between the narrow bands. The cones are protected by $\mathcal{C}_2\mathcal{T}$ which is preserved even in the presence of the electric field [61, 62]. The presence of gap between narrow bands in TBG with hBN is because the substrate breaks \mathcal{C}_2 . To introduce an electric field in our TB model, we add an on-site potential with different sign on each graphene layer. The electrostatic potential is given by $\Delta V_i = z_i \mathcal{E} \cdot \mathbf{e}_z$ being \mathcal{E} the vertical electric field which has positive value if its direction is the same to \mathbf{e}_z (upward) and z_i is the z coordinate of atom i . We choose the reference plane for electrostatic potential at the midpoint of two graphene layers. Then, an extra onsite term is added in the Hamiltonian in Eq. (5) as

$$H_V = \Delta V_i |\mathbf{R}_i\rangle \langle \mathbf{R}_i|. \quad (12)$$

This onsite term is positive in the upper graphene plane and negative in the lower graphene plane for electric field with upward direction. We neglect this onsite potential in the N and B atoms since the onsite energies of N and B are much larger than this extra onsite term induced by the external electric field.

In pristine TBG, mirror symmetry ensures the degeneracy of the Dirac cones. In the presence of the electric field this symmetry is broken because at each layer an onsite potential with different sign is introduced and this is the main source of the gap between narrow bands (see Ref. [12] for further details). Indeed, this effect can be clearly seen in Fig. 3. In the TBG/hBN system (top row of Fig. 3) the degeneracy of the Dirac cones (or layer degeneracy) is broken because the hBN is acting on a

single graphene layer. Notice that the energy bands corresponding to each valley are shifted in opposite directions. The effect of the electric field is a further shifting of the energy bands and this effectively looks like a gap closing as shown in the DOS in Fig. 3(b) and (e). For the encapsulated system (bottom row of Fig. 3) we can see that the electric field shifts the energy bands, however, the effect of the field is less drastic because the layer degeneracy is preserved (see Fig. 2(b)). The effects of the layer degeneracy breaking can be observed in the density maps in Fig. 3(c) and (f). As expected in TBG systems, the charge is strongly localized around the AA centers even with an hBN substrate. In the presence of an electric field in Fig. 3, the charge is localized around the ACB region for both the C3 and the C4 bands, whereas at zero field in Fig. 1, the charge is localized in different regions for different bands. This difference is originated from the layer degeneracy breaking. There is a charge transfer between two graphene layers. The induced electric field enhances the interlayer interaction by producing an energy difference between two graphene layers, thus re-distributing the charge density. In the presence of both hBN and an electric field, the charge is unevenly distributed between layers inducing polarization. The latter is stronger in the suspended situation, as shown in Fig. 3(c). We have found that the layer polarization is a physical phenomena resulting from layer degeneracy breaking. While this degeneracy can be recovered in TBG by encapsulation with hBN, c.f Fig. 2, with an electric field is always broken. Our results indicate that in TBG samples suspended or encapsulated with hBN, the presence of an electric field polarizes the

charges within the layers. From time-reversal symmetry, the band energy shift is opposite in each valley and therefore the charge polarization is a valley phenomena. The total charge distribution in each layer is the same if both valleys are considered. This effect can be used to tune the strength of the valley Hall effect or the anomalous Hall effect in the presence of a magnetic field.

V. STACKING DEPENDENT BAND TOPOLOGY

Theoretical studies analyzing the topology and correlated effects in the presence of a substrate usually introduce the effect of the hBN by considering only the mass term [21, 63, 64]. While this approach is valid, in experiments the TBG samples are supported on [65] or encapsulated [66, 67] with hBN and this may result in different band topology. In this section, we show that the topological properties of TBG/hBN and hBN/TBG/hBN can be completely different. In both systems, the breaking of inversion symmetry in the TB model allows for a non-zero Berry curvature with opposite signs in each valley. Because of time-reversal symmetry, the Berry curvature in each valley has opposite sign and hence the total Chern number of a given band is zero. However, the topological invariants can be defined for each valley [31, 68]. The finite Berry curvature for a single valley is given by

$$\Omega_{\vec{k},l} = 2 \operatorname{Im} \left\{ \left\langle \partial_{k_x} \psi_{\vec{k},l} \middle| \partial_{k_y} \psi_{\vec{k},l} \right\rangle \middle| \partial_{k_x} \psi_{\vec{k},l} \right\rangle \middle| \partial_{k_y} \psi_{\vec{k},l} \right\rangle \right\}, \quad (13)$$

where l is a band index with energy $E_l(\vec{k})$, momentum $\vec{k} = \{k_x, k_y\}$ and wavefunctions $\psi_{\vec{k},l}$. For the different stacking configurations, the bands for each valley are isolated and their Berry curvature is well-defined. Therefore, we can assign a valley Chern number, \mathcal{C}_l , to the band l which is given by the integral of the Berry curvature in the moiré Brillouin zone

$$\mathcal{C}_l = \frac{1}{2\pi} \int_{\text{mBZ}} d^2\vec{k} \Omega_{\vec{k},l}. \quad (14)$$

We use the algorithm in Ref. [69] and the continuum model in Ref. [12] to compute the Berry curvature. To simplify our analysis we are considering only the mass terms, here, $\Delta_{b/t}$ are mass terms acting in the bottom and top graphene layer, respectively, mimicking an hBN substrate, we also consider the bands only at charge neutrality because the band topology may also depend on the filling fraction [70]. Figure 4 displays the set of valley Chern numbers for the systems considered in this work. We can distinguish three topological phases with Chern numbers for each band, $\{\mathcal{C}_b, \mathcal{C}_t\}$, identified as $P_1 = \{-1, 1\}$, $P_2 = \{0, 0\}$ and $P_3 = \{1, -1\}$. The phases P_1 and P_3 are the usual phases found in the presence of a single hBN layer. However, for the encapsulated situation there are different possibilities, such as the case of Fig. 2(b) where the corresponding phase is P_3 because

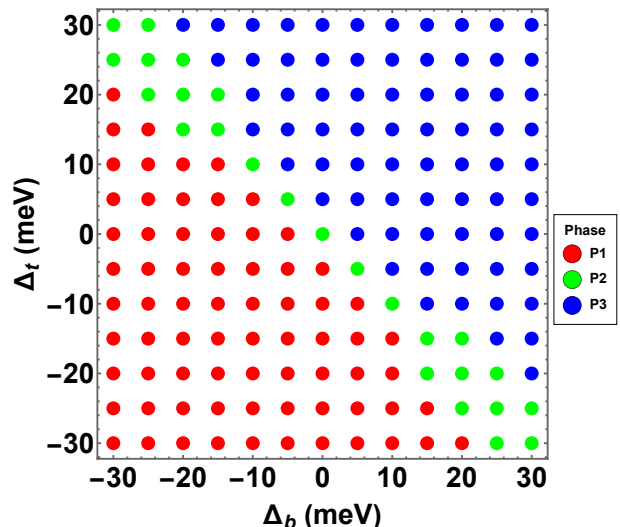


FIG. 4. Topological phases of the two narrow middle bands of TBG with hBN as a function of the mass terms. We distinguish three topological phases, $\{\mathcal{C}_b, \mathcal{C}_t\}$, identified as $P_1 = \{-1, 1\}$, $P_2 = \{0, 0\}$ and $P_3 = \{1, -1\}$

we found that the mass terms have the same sign and magnitude (~ 30 meV). Additional stacking configurations may give different topological phases, for example, if we swap the Boron by the Nitrogen in the stack in Fig. 2(a) the mass terms change sign and the phase can be P_1 or P_2 . Interestingly, in the encapsulated situation, if each hBN is nearly aligned with their adjacent TBG the topological phase is P_2 if the mass terms have different sign and similar magnitude. This situation can be achieved, for example, by aligning both hBN layers with respect to each other or by rotating them in opposite directions with the same angle. The dominant contribution in TBG/hBN structures are the mass terms. If the angle between hBN and its adjacent graphene layer is increased, these terms decrease and may survive up to 3° of alignment (see Ref. [12] for further details). Therefore, Fig. 4 is a general result and indicates the rich band topology of suspended and encapsulated samples of TBG with hBN that can be obtained near alignment.

VI. CONCLUSIONS

We have studied the effects of an hexagonal boron nitride substrate on the electronic and topological properties of TBG. In particular, we have studied TBG on hBN and TBG encapsulated with hBN. By using a real space tight-binding method in combination with semi-classical molecular dynamics, we calculate the band structure, DOS, LDOS and analyze the stacking-dependent topological properties of these systems. We find that the substrate significantly modifies the electronic properties of the TBG due to the broken \mathcal{C}_2 symmetry. Compared to the free standing TBG, the narrow bands have been

strongly modified and separated by a gap of around 30 meV. In the TBG/hBN system we found that the substrate induces a layer degeneracy breaking which results in a splitting of the TB band structure and uneven distribution of the LDOS in a single layer. In the encapsulated TBG/hBN system we found that the layer degeneracy is recovered if the twist angle between each graphene and its nearest hBN layer is of the same magnitude. Both considered heterostructures are also strongly sensitive to a perpendicular electric field. A direct consequence of this field is the shifting of the energy bands which may look as a gap closing in the DOS. Interestingly, by mapping the real space distribution of the wavefunctions we found that the electric field polarizes the charge in each valley which can be used to tune the strength of the valley Hall effect or the anomalous Hall effect in the presence of a magnetic field. Finally, by calculating the valley Chern numbers, we found that depending on the induced mass gap, different topological phases can be obtained. Because the mass gap is a consequence of the degree of the hBN alignment, this result may give a simple methodology to identify the encapsulating conditions in TBG/hBN samples.

ACKNOWLEDGMENTS

This work was supported by the National Natural Science Foundation of China (Grant No. 11774269) and the National Key R&D Program of China (Grant No. 2018YFA0305800). IMDEA Nanociencia acknowledges support from the “Severo Ochoa” Programme for Centres of Excellence in R&D (CEX2020-001039-

S / AEI / 10.13039/501100011033). Z.Z. acknowledges support funding from the European Union’s Horizon 2020 research and innovation programme under the Marie Skłodowska-Curie grant agreement No 101034431 and from the “Severo Ochoa” Programme for Centres of Excellence in R&D (CEX2020-001039-S / AEI / 10.13039/501100011033). P.A.P and F.G. acknowledge funding from the European Commission, within the Graphene Flagship, Core 3, grant number 881603 and from grants NEMAT2D (Comunidad de Madrid, Spain), SprQuMat (Ministerio de Ciencia e Innovación, Spain). Numerical calculations presented in this paper have been performed on the supercomputing system in the Supercomputing Center of Wuhan University.

Appendix A: The electronic properties of free standing twisted bilayer graphene

In this section, we briefly describe the properties of pristine TBG. The hopping parameters in Eq. (10) are given such that the “magic” angle is at 1.21° . Figure 5(a) show the band structure of TBG with a twist angle $\theta_{tbg} = 1.05^\circ$. The DOS display the van Hove singularities (vHs) due to the two narrow bands. Figure 5(b) shows that the contribution to each sublattice to the LDOS is identical due to the preserved sublattice, layer and valley symmetries. In Fig. 5(c) we show the LDOS maps in real space for energies at the vHs. They have the familiar “fidget-spinner” shape [62] with the states localized around the AA region.

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- [1] Y. Cao, V. Fatemi, A. Demir, S. Fang, S. L. Tomarken, J. Y. Luo, J. D. Sanchez-Yamagishi, K. Watanabe, T. Taniguchi, E. Kaxiras, R. C. Ashoori, and P. Jarillo-Herrero, *Nature* **556**, 80 (2018).
 - [2] Y. Cao, V. Fatemi, S. Fang, K. Watanabe, T. Taniguchi, E. Kaxiras, and P. Jarillo-Herrero, *Nature* **556**, 43 (2018).
 - [3] S. Carr, D. Massatt, S. Fang, P. Cazeaux, M. Luskin, and E. Kaxiras, *Phys. Rev. B* **95** (2017).
 - [4] R. Bistritzer and A. H. MacDonald, *PNAS* **108**, 12233 (2011).
 - [5] Z. Zhang, Y. Wang, K. Watanabe, T. Taniguchi, K. Ueno, E. Tutuc, and B. J. LeRoy, *Nat. Phys.* **16**, 1093 (2020).
 - [6] Z. Zhan, Y. Zhang, P. Lv, H. Zhong, G. Yu, F. Guinea, J. Á. Silva-Guillén, and S. Yuan, *Phys. Rev. B* **102** (2020).
 - [7] D. M. Kennes, L. Xian, M. Claassen, and A. Rubio, *Nat. Commun.* **11** (2020).
 - [8] C. R. Woods, P. Ares, H. Nevison-Andrews, M. J. Holwill, R. Fabregas, F. Guinea, A. K. Geim, K. S. Novoselov, N. R. Walet, and L. Fumagalli, *Nat. Commun.* **12** (2021).
 - [9] N. R. Walet and F. Guinea, *Phys. Rev. B* **103** (2021).
 - [10] T. Cao, Z. Li, D. Y. Qiu, and S. G. Louie, *Nano Lett.* **16**, 5542 (2016).
 - [11] L. Wang, S. Zihlmann, M.-H. Liu, P. Makk, K. Watanabe, T. Taniguchi, A. Baumgartner, and C. Schönenberger, *Nano Lett.* **19**, 2371 (2019).
 - [12] M. Long, P. A. Pantaleón, Z. Zhan, F. Guinea, J. Á. Silva-Guillén, and S. Yuan, *Npj Comput. Mater.* **8** (2022).
 - [13] J. Shi, J. Zhu, and A. H. MacDonald, *Phys. Rev. B* **103**, 075122 (2021).
 - [14] T. Cea, P. A. Pantaleón, and F. Guinea, *Phys. Rev. B* **102**, 155136 (2020).
 - [15] J. Shin, Y. Park, B. L. Chittari, J.-H. Sun, and J. Jung, *Phys. Rev. B* **103**, 075423 (2021).
 - [16] D. Mao and T. Senthil, *Phys. Rev. B* **103**, 115110 (2021).
 - [17] T. Cea, P. A. Pantaleón, and F. Guinea, *Phys. Rev. B* **102**, 155136 (2020).
 - [18] B. L. Chittari, G. Chen, Y. Zhang, F. Wang, and J. Jung, *Phys. Rev. Lett.* **122**, 016401 (2019).
 - [19] F. Mukai, K. Horii, R. Ebisuoka, K. Watanabe, T. Taniguchi, and R. Yagi, *Commun. Phys.* **4** (2021).
 - [20] J. Shin, Y. Park, B. L. Chittari, J.-H. Sun, and J. Jung, *Phys. Rev. B* **103**, 075423 (2021).

- [21] Y.-H. Zhang, D. Mao, and T. Senthil, *Phys. Rev. Res.* **1** (2019).
- [22] S. Plimpton, *J. Comput. Phys.* **117**, 1 (1995).
- [23] D. W. Brenner, O. A. Shenderova, J. A. Harrison, S. J. Stuart, B. Ni, and S. B. Sinnott, *J. Phys. Condens. Matter* **14**, 783 (2002).
- [24] A. N. Kolmogorov and V. H. Crespi, *Phys. Rev. B* **71**, 235415 (2005).
- [25] G. J. Slotman, M. M. van Wijk, P.-L. Zhao, A. Fasolino, M. I. Katsnelson, and S. Yuan, *Phys. Rev. Lett.* **115**, 186801 (2015).
- [26] P. Moon and M. Koshino, *Phys. Rev. B* **90** (2014).
- [27] P. San-Jose, A. Gutiérrez-Rubio, M. Sturla, and F. Guinea, *Phys. Rev. B* **90**, 075428 (2014).
- [28] J. Jung, E. Laksono, A. M. Dasilva, A. H. Macdonald, M. Mucha-Kruczyński, and S. Adam, *Phys. Rev. B* **96**, 085442 (2017).
- [29] B. Sachs, T. O. Wehling, M. I. Katsnelson, and A. I. Lichtenstein, *Phys. Rev. B* **84**, 195414 (2011).
- [30] M. Kindermann, B. Uchoa, and D. L. Miller, *Phys. Rev. B* **86**, 115415 (2012).
- [31] P. San-Jose, A. Gutiérrez-Rubio, M. Sturla, and F. Guinea, *Phys. Rev. B* **90**, 115152 (2014).
- [32] J. R. Wallbank, A. A. Patel, M. Mucha-Kruczyński, A. K. Geim, and V. I. Fal'ko, *Phys. Rev. B* **87**, 245408 (2013).
- [33] X. Lin and J. Ni, *Phys. Rev. B* **100**, 195413 (2019).
- [34] H. Ochoa, E. V. Castro, M. I. Katsnelson, and F. Guinea, *Phys. Rev. B* **83** (2011).
- [35] S. Yuan, H. D. Raedt, and M. I. Katsnelson, *Phys. Rev. B* **82** (2010).
- [36] Y. Li, Z. Zhan, X. Kuang, Y. Li, and S. Yuan, *Comput. Phys. Commun.* **285**, 108632 (2023).
- [37] TBPLaS, <http://www.tbplas.net/index.html> (2022).
- [38] R. Haydock, V. Heine, and M. Kelly, *J. Phys. C: Solid State Phys.* **5**, 2845 (1972).
- [39] H. Shi, Z. Zhan, Z. Qi, K. Huang, E. van Veen, J. Á. Silva-Guillén, R. Zhang, P. Li, K. Xie, H. Ji, M. I. Katsnelson, S. Yuan, S. Qin, and Z. Zhang, *Nat. Commun.* **11**, 1 (2020).
- [40] J. Jung, A. M. DaSilva, A. H. MacDonald, and S. Adam, *Nat. Commun.* **6** (2015).
- [41] B. Hunt, J. D. Sanchez-Yamagishi, A. F. Young, M. Yankowitz, B. J. LeRoy, K. Watanabe, T. Taniguchi, P. Moon, M. Koshino, P. Jarillo-Herrero, and R. C. Ashoori, *Science* **340**, 1427 (2013).
- [42] J. C. W. Song, A. V. Shytov, and L. S. Levitov, *Phys. Rev. Lett.* **111**, 266801 (2013).
- [43] F. Amet, J. R. Williams, K. Watanabe, T. Taniguchi, and D. Goldhaber-Gordon, *Phys. Rev. Lett.* **110**, 216601 (2013).
- [44] R. V. Gorbachev, J. C. W. Song, G. L. Yu, A. V. Kretinin, F. Withers, Y. Cao, A. Mishchenko, I. V. Grigorieva, K. S. Novoselov, L. S. Levitov, and A. K. Geim, *Science* **346**, 448 (2014).
- [45] Z.-G. Chen, Z. Shi, W. Yang, X. Lu, Y. Lai, H. Yan, F. Wang, G. Zhang, and Z. Li, *Nat. Commun.* **5**, 4461 (2014).
- [46] M. Yankowitz, J. Xue, and B. J. LeRoy, *J. Phys. Condens. Matter* **26**, 303201 (2014).
- [47] D. Wong, Y. Wang, J. Jung, S. Pezzini, A. M. DaSilva, H.-Z. Tsai, H. S. Jung, R. Khajeh, Y. Kim, J. Lee, S. Kahn, S. Tollabimazraehno, H. Rasool, K. Watanabe, T. Taniguchi, A. Zettl, S. Adam, A. H. MacDonald, and M. F. Crommie, *Phys. Rev. B* **92**, 155409 (2015).
- [48] M. Lee, J. R. Wallbank, P. Gallagher, K. Watanabe, T. Taniguchi, V. I. Fal'ko, and D. Goldhaber-Gordon, *Science* **353**, 1526 (2016).
- [49] E. Wang, X. Lu, S. Ding, W. Yao, M. Yan, G. Wan, K. Deng, S. Wang, G. Chen, L. Ma, J. Jung, A. V. Fedorov, Y. Zhang, G. Zhang, and S. Zhou, *Nat. Phys.* **12**, 1111 (2016).
- [50] M. Yankowitz, J. Jung, E. Laksono, N. Leconte, B. L. Chittari, K. Watanabe, T. Taniguchi, S. Adam, D. Graf, and C. R. Dean, *Nature* **557**, 404 (2018).
- [51] H. Kim, N. Leconte, B. L. Chittari, K. Watanabe, T. Taniguchi, A. H. MacDonald, J. Jung, and S. Jung, *Nano Lett.* **18**, 7732 (2018).
- [52] P. A. Pantaleón, T. Cea, R. Brown, N. R. Walet, and F. Guinea, *2D Mater.* **8**, 044006 (2021).
- [53] S. Moriyama, Y. Morita, K. Komatsu, K. Endo, T. Iwasaki, S. Nakaharai, Y. Noguchi, Y. Wakayama, E. Watanabe, D. Tsuya, *et al.*, [arXiv preprint arXiv:1901.09356](https://arxiv.org/abs/1901.09356) (2019).
- [54] G. Chen, A. L. Sharpe, P. Gallagher, I. T. Rosen, E. J. Fox, L. Jiang, B. Lyu, H. Li, K. Watanabe, T. Taniguchi, J. Jung, Z. Shi, D. Goldhaber-Gordon, Y. Zhang, and F. Wang, *Nature* **572**, 215 (2019).
- [55] M. He, Y. Li, J. Cai, Y. Liu, K. Watanabe, T. Taniguchi, X. Xu, and M. Yankowitz, *Nat. Phys.* **17**, 26 (2020).
- [56] S. Carr, C. Li, Z. Zhu, E. Kaxiras, S. Sachdev, and A. Kruchkov, *Nano Lett.* **20**, 3030 (2020).
- [57] C. Lei, L. Linhart, W. Qin, F. Libisch, and A. H. MacDonald, *Phys. Rev. B* **104**, 035139 (2021).
- [58] A. Lopez-Bezanilla and J. L. Lado, *Phys. Rev. Res.* **2** (2020).
- [59] B. Tsim, N. N. T. Nam, and M. Koshino, *Phys. Rev. B* **101**, 125409 (2020).
- [60] P. San-Jose and E. Prada, *Phys. Rev. B* **88**, 121408 (2013).
- [61] P. Moon, Y.-W. Son, and M. Koshino, *Phys. Rev. B* **90**, 155427 (2014).
- [62] H. C. Po, L. Zou, A. Vishwanath, and T. Senthil, *Phys. Rev. X* **8** (2018).
- [63] C. Repellin and T. Senthil, *Phys. Rev. Res.* **2**, 023238 (2020).
- [64] N. Bultinck, S. Chatterjee, and M. P. Zaletel, *Phys. Rev. Lett.* **124**, 166601 (2020).
- [65] M. Serlin, C. L. Tschirhart, H. Polshyn, Y. Zhang, J. Zhu, K. Watanabe, T. Taniguchi, L. Balents, and A. F. Young, *Science* **367**, 900 (2019).
- [66] A. L. Sharpe, E. J. Fox, A. W. Barnard, J. Finney, K. Watanabe, T. Taniguchi, M. A. Kastner, and D. Goldhaber-Gordon, *Science* **365**, 605 (2019).
- [67] C. Ma, Q. Wang, S. Mills, X. Chen, B. Deng, S. Yuan, C. Li, K. Watanabe, T. Taniguchi, X. Du, F. Zhang, and F. Xia, *Nano Lett.* **20**, 6076 (2020).
- [68] J. C. Song, P. Samutpraphoot, and L. S. Levitov, *PNAS* **112**, 10879 (2015).
- [69] T. Fukui, Y. Hatsugai, and H. Suzuki, *JPSJ* **74**, 1674 (2005).
- [70] P. A. Pantaleón, V. T. Phong, G. G. Naumis, and F. Guinea, *Phys. Rev. B* **106** (2022).

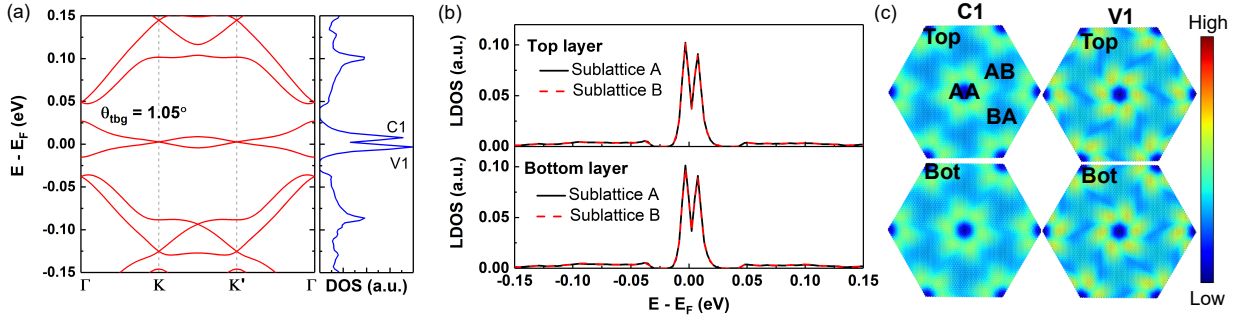


FIG. 5. (a) Band structure and density of states, (b) LDOS of the sublattice A and B in the AA stacking region and (c) LDOS mapping in real space for free standing twisted bilayer graphene with $\theta = 1.05^\circ$. The LDOS mapping are obtained via the TBPM method. The corresponding energies of the eigenstates are illustrated in the DOS.